

# MF6/MT18 format status

$P(\nu)$ ,  $P(\nu_g)$  and multiplicity-dependent spectra

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CSEWG, November 2017

# MF6 MT18: including P(nu) and P(nug)

```
[MAT, 6, MT/ ZA, AWR,   JP, LCT,  NK,   0]HEAD  
      <TAB1 and LAW-dependent structure for product 1>  
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```

```
      <repeat TAB1 and LAW-dependent structures>  
      <for the rest of the NK subsections      >  
-----
```

```
[MAT, 6, MT/ 0.0, 0.0,   0,   0,   0,   0]SEND
```

```
[MAT, 6, MT/ ZAP, AWP, LIP, LAW,  NR,  NP/ Eint / yi(E)]TAB1  
      <LAW-dependent structure for product 1>
```

- ❖ NK is the sum of neutron and gamma sections (numax+1+nugmax+1)
- ❖ JP=10\*JPP+JPN
- ❖ ZAP and AWP is either 1 for neutrons or 0 for gammas
- ❖ Neutrons are listed first
- ❖ LIP used to identify the type of spectrum. Suggested to identify multiplicity

# JP flag

**JPX=0** This option indicates that the spectrum averaged over multiple outgoing particles and the average multiplicity is given. In this case, the usual MF=6 interpretation of  $y_i$  and  $f_i$  would hold. As this is fission, the multiplicity is already given in MF=1/MT=456 for neutrons and MF=18/MT=12 for photons and the outgoing distributions in MT=18/MF=4 & 5 for neutrons or MT=18/MF=14 & 15 for photons. Allowing either neutron or photon datasets here would require duplicating data, potentially creating synchronization errors. Therefore, NKN=0 no neutron data given and if NKP=0 no photon data is given. Indeed, this entire file (MT=18/MF=6) can be omitted if NK=NKP+NKN=0.

**JPX=1** This option indicates that the average spectrum is given as well as the probability functions  $P_i(\nu_i, E)$  for particle  $i$ . In this case,  $y_i$  stores the probability function  $P_i(\nu_i, E)$ , but the usual MF=6 interpretation of  $f_i$  holds. To use this option,

- Set NKN =  $(\nu_{n,\max} + 1)$  for  $JPN \geq 1$  &  $JPP=0$ , NKP =  $(\nu_{\gamma,\max} + 1)$  for  $JPP \geq 1$  &  $JPN=0$  and NK=NKN+NKP =  $(\nu_{n,\max} + 1) + (\nu_{\gamma,\max} + 1)$  for  $JPN, JPP \geq 1$ .
- For  $\nu_i = 0$ , set product 1 as follows: LIP=0, LAW=0 (unspecified) and store  $P_i(\nu_i = 0, E)$  in  $y_i$ .
- For each additional  $\nu_i$ , use LIP=0 and LAW<0 and  $P_i(\nu_i, E)$  for  $y_i$ . LAW<0 signals to the processing codes to look to MF=4 & 5 for the neutron's outgoing energy spectrum and angular distribution and to MF=14 & 15 for gamma's outgoing energy spectrum and angular distribution.

**JPX=2** This indicates that the probability functions  $P_i(\nu_i, E)$  and the spectra  $\bar{\chi}_i(\nu_i, E, E'_i)$  are given. In this case,  $y_i$  stores the probability function  $P_i(\nu_i, E)$ , and  $f_i$  stores the spectrum  $\bar{\chi}_i(\nu_i, E, E'_i)$  for the  $\nu_i$  particles of type  $i$ . To use this option,

- Set NKN or NKP as above for JPX=1.
- For  $\nu_i = 0$ , set product 1 as follows: LIP=0, LAW=0 (unspecified) and store  $P_i(\nu_i = 0, E)$  in  $y_i$ . As there are no particles emitted in this case, the emitted spectrum is undefined.
- For each additional  $\nu_i$ , use LIP=0 and LAW>0 and  $P_i(\nu_i, E)$  for  $y_i$ . Use  $\bar{\chi}_i(\nu_i, E, E'_i)$  as the  $f_i$  for this “product” and represent it with the appropriate LAW.

# Prompt fission spectrum decomposition

$$\sum_{\nu=0}^{\nu_{max}} P(\nu, E) = 1$$

$$\sum_{\nu=0}^{\nu_{max}} \nu P(\nu, E) = \bar{\nu}$$

$$\bar{\chi}(E, E') = \frac{1}{\bar{\nu}} \sum_{\nu=1}^{\nu_{max}} P(\nu, E) \sum_{k=1}^{\nu} \chi_{\nu}(k, E, E') = \frac{1}{\bar{\nu}} \sum_{\nu=1}^{\nu_{max}} P(\nu, E) \nu \bar{\chi}_{\nu}(E, E')$$

**Issue with interpolation: consistency with 1/456 and/or 12/18**

A. Trkov:

$$\tilde{P}(\nu, E) = \frac{\nu}{\bar{\nu}} P(\nu, E)$$

P. Talou:

Save parameters of a function  
(new format/too late)

# Example: JPN=1 and JPP=1

9.223500+4 2.330248+2	11	<del>0</del>	56	09228 6 18
1.000000+0 1.000000+0	0	<del>0</del>	1	229228 6 18
	22	2		9228 6 18
1.000000-5 3.817884-2 5.000000+5 3.428548-2 1.000000+6 3.116930-29228 6 18				
1.500000+6 2.761014-2 2.000000+6 2.434288-2 2.500000+6 2.134849-29228 6 18				
3.000000+6 1.854712-2 4.000000+6 1.390180-2 5.000000+6 9.546429-39228 6 18				
6.000000+6 6.033709-3 7.000000+6 3.660794-3 8.000000+6 2.471095-39228 6 18				
9.000000+6 1.613662-3 1.000000+7 1.051794-3 1.100000+7 6.837755-49228 6 18				
1.200000+7 4.405975-4 1.300000+7 2.744386-4 1.400000+7 1.639486-49228 6 18				
1.500000+7 9.544946-5 2.000000+7 7.144631-6 2.500000+7 6.927065-79228 6 18				
3.000000+7 5.338292-8				9228 6 18
1.000000+0 1.000000+0	0	-5	1	229228 6 18
	22	2		9228 6 18
1.000000-5 1.605149-1 5.000000+5 1.510252-1 1.000000+6 1.429648-19228 6 18				
1.500000+6 1.331807-1 2.000000+6 1.235745-1 2.500000+6 1.141573-19228 6 18				
3.000000+6 1.047219-1 4.000000+6 8.740858-2 5.000000+6 6.858005-29228 6 18				
6.000000+6 5.053576-2 7.000000+6 3.588973-2 8.000000+6 2.725382-29228 6 18				
9.000000+6 2.011503-2 1.000000+7 1.475440-2 1.100000+7 1.075457-29228 6 18				
1.200000+7 7.756523-3 1.300000+7 5.432451-3 1.400000+7 3.670901-39228 6 18				
1.500000+7 2.421633-3 2.000000+7 3.135936-4 2.500000+7 4.717257-59228 6 18				
3.000000+7 5.643180-6				9228 6 18

LAW<0: look for the spectrum in MF5 (neutrons) and MF15 (photons)

# Example JPN=1 and JPP=2

0.000000+0	0.000000+0	0	0	1	229228	6	18	
22	2				9228	6	18	
1.000000-5	1.057797-3	1.000000+6	2.288868-4	2.000000+6	0.000000+0	9228	6	18
3.000000+6	3.038176-4	4.000000+6	1.801776-4	5.000000+6	0.000000+0	9228	6	18
6.000000+6	4.980726-5	7.000000+6	7.701545-5	8.000000+6	0.000000+0	9228	6	18
9.000000+6	6.705273-5	1.000000+7	5.606544-5	1.100000+7	0.000000+0	9228	6	18
1.200000+7	3.666197-5	1.300000+7	2.744603-5	1.400000+7	0.000000+0	9228	6	18
1.500000+7	1.286589-5	1.600000+7	1.106171-5	1.700000+7	0.000000+0	9228	6	18
1.800000+7	9.187529-6	1.900000+7	8.052606-6	2.000000+7	0.000000+0	9228	6	18
3.000000+7	5.650443-7				9228	6	18	
0.000000+0	0.000000+0	0	1	1	229228	6	18	
22	2				9228	6	18	
1.000000-5	5.879624-3	1.000000+6	1.719723-3	2.000000+6	0.000000+0	9228	6	18
3.000000+6	2.025959-3	4.000000+6	1.288339-3	5.000000+6	0.000000+0	9228	6	18
6.000000+6	4.221620-4	7.000000+6	6.287373-4	8.000000+6	0.000000+0	9228	6	18
9.000000+6	5.540422-4	1.000000+7	4.732953-4	1.100000+7	0.000000+0	9228	6	18
1.200000+7	3.267141-4	1.300000+7	2.546359-4	1.400000+7	0.000000+0	9228	6	18
1.500000+7	1.329930-4	1.600000+7	1.158163-4	1.700000+7	0.000000+0	9228	6	18
1.800000+7	9.674664-5	1.900000+7	8.559761-5	2.000000+7	0.000000+0	9228	6	18
3.000000+7	7.863344-6				9228	6	18	
0.000000+0	0.000000+0	1	2	1	29228	6	18	
2	2				9228	6	18	
0.000000+0	1.000000-5	0	0	592	2969228	6	18	
1.000000-5	0.000000+0	1.000000+4	1.107088-7	3.000000+4	3.321263-79228	6	18	
5.000000+4	5.535439-7	7.000000+4	7.749624-7	9.000000+4	9.963810-79228	6	18	
1.100000+5	1.217795-6	1.300000+5	1.290769-6	1.500000+5	1.398886-69228	6	18	
1.700000+5	1.049633-6	1.900000+5	1.150249-6	2.100000+5	1.473374-69228	6	18	
2.300000+5	8.199414-7	2.500000+5	8.630149-7	2.700000+5	7.311250-79228	6	18	
2.900000+5	1.019560-6	3.100000+5	7.560036-7	3.300000+5	8.683659-79228	6	18	

# Life is complicated

The evaluation is a mix of experiment and theoretical simulations, and multiplicity-dependent spectra come from theory only => need to be adjusted

